

**Numerical Simulation of Heat Transfer in Materials
with Anisotropic Thermal Conductivity:
A Finite Volume Scheme to Handle Complex Geometries**

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Complex Sample Domain from Crystal Growth

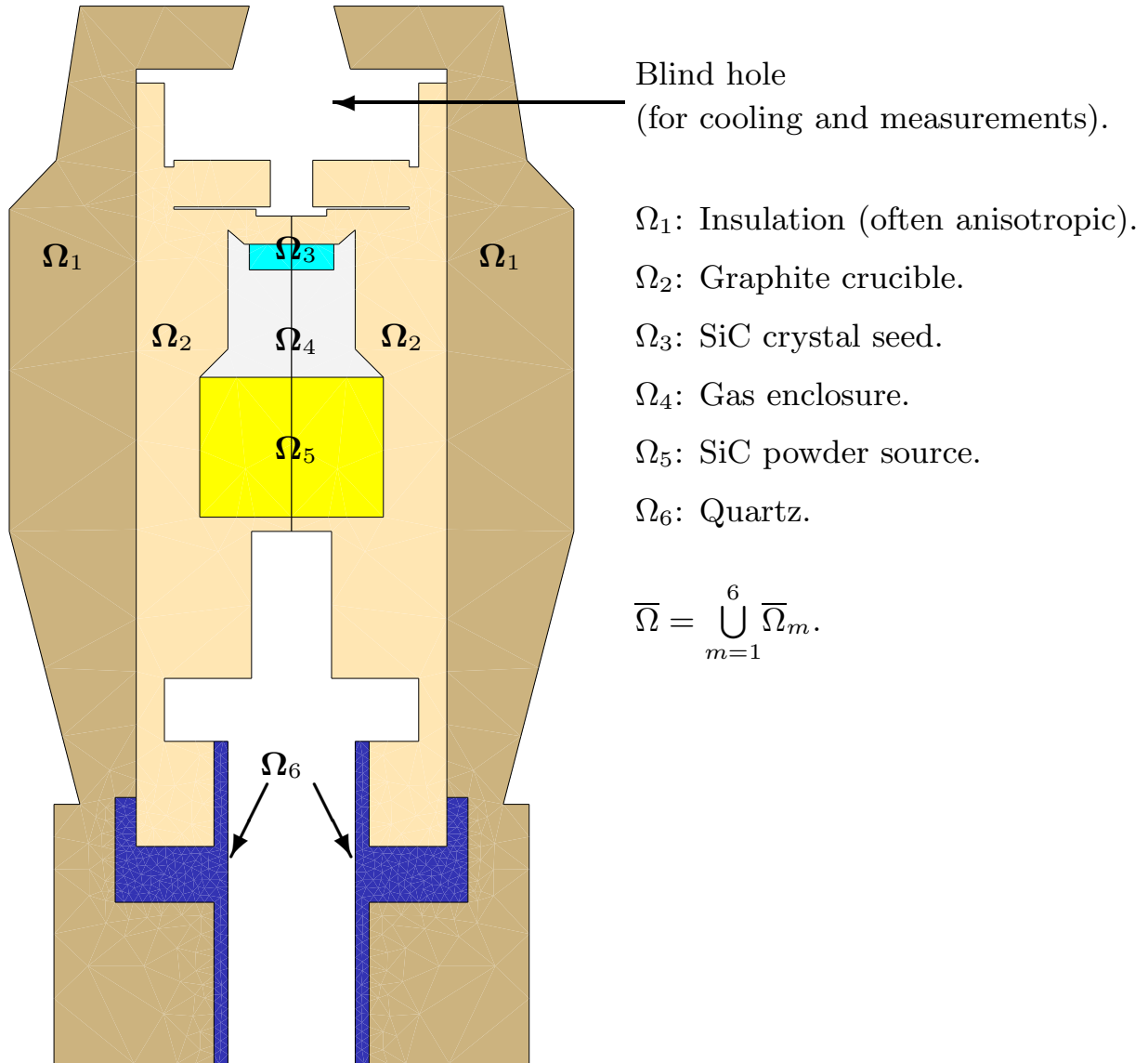


Figure 1: Axisymmetric domain representing a growth apparatus used in silicon carbide single crystal growth by physical vapor transport (PVT). The geometry is a modified version of K. Semmelroth et al., J. Phys.-Condes. Matter 16 (2004).

Model for Stationary Anisotropic Heat Conduction

$$-\operatorname{div}(K_m(\theta) \nabla \theta) = f_m \quad \text{in } \Omega_m \quad (m \in M),$$

θ : absolute temperature, K_m : symmetric and positive definite tensor of thermal conductivity, f_m : heat sources, Ω_m : domain of material m .

Assumed form of K_m :

$$K_m(\theta) = \kappa_{i,j}^m(\theta), \quad \text{where } \kappa_{i,j}^m(\theta) = \begin{cases} \alpha_i^m \kappa_{\text{iso}}^m(\theta) & \text{for } i = j, \\ 0 & \text{for } i \neq j. \end{cases}$$

Interface Conditions on $\overline{\Omega}_{m_1} \cap \overline{\Omega}_{m_2}$:

$$K_{m_1}(\theta) \nabla \theta \upharpoonright_{\overline{\Omega}_{m_1}} \bullet \mathbf{n}_{m_1} = K_{m_2}(\theta) \nabla \theta \upharpoonright_{\overline{\Omega}_{m_2}} \bullet \mathbf{n}_{m_1}.$$

\upharpoonright : restriction, \mathbf{n}_{m_1} : outer unit normal vector to material m_1 .

Boundary Conditions: Dirichlet, Robin

$$\theta = \theta_{\text{Dir}} \quad \text{on } \overline{\Gamma}_{\text{Dir}},$$

$$-K_m(\theta) \nabla \theta \bullet \mathbf{n}_m = \xi_m (\theta - \theta_{\text{ext}}) \quad \text{on } \Gamma_{\text{Rob}} \cap \partial\Omega_m, m \in M,$$

Finite Volume Discretization

$\Sigma_m = (\sigma_{m,i})_{i \in I_m}$ conforming triangulation of Ω_m satisfying the **constrained Delaunay property**: If γ is an interior edge of Σ_m , α and β the angles opposite to γ , then $\alpha + \beta \leq \pi$. If $\gamma \subseteq \partial\Omega_m$ is a boundary edge of Σ_m , α the angle opposite γ , then $\alpha \leq \pi/2$.

$V(\sigma_{m,i}) = v_{i,j}^m : j \in \{1, 2, 3\}$: Set of vertices.

$V := \bigcup_{m \in M, i \in I_m} V(\sigma_{m,i})$.

$\omega_v := \{x \in \Omega : \|x - v\|_2 < \|x - z\|_2 \text{ for each } z \in V \setminus \{v\}\}$,

$\omega_{m,v} := \omega_v \cap \Omega_m, \quad V_m := \{z \in V : \omega_{m,z} \neq \emptyset\}$.

$$A_m = (a_{i,j}^m), \quad a_{i,j}^m := \begin{cases} \alpha_i^m & \text{for } i = j, \\ 0 & \text{for } i \neq j. \end{cases}$$

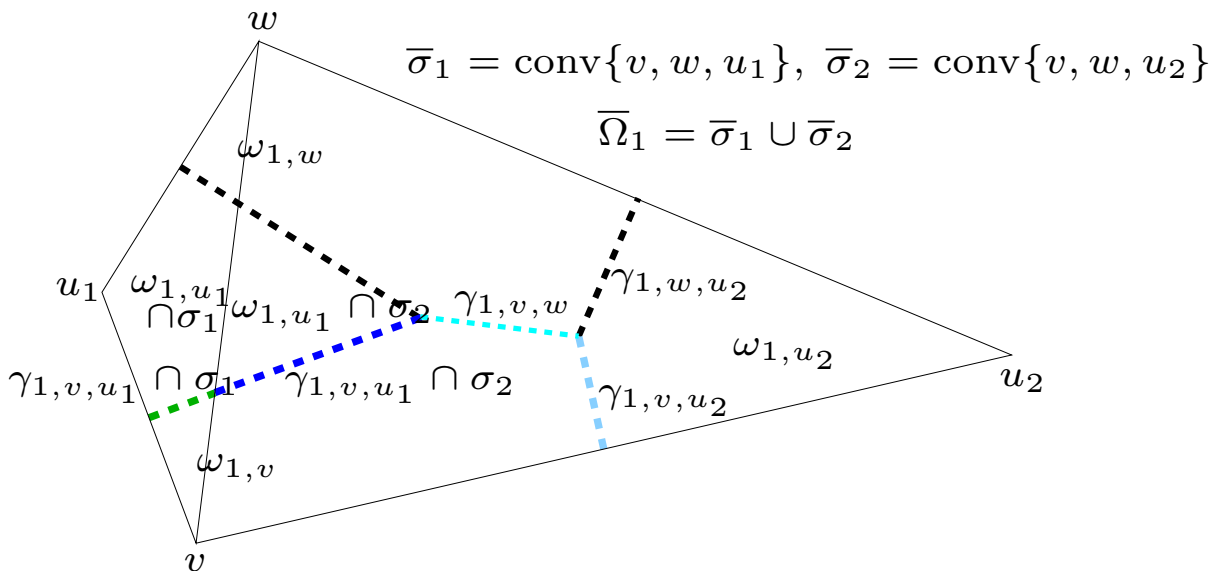


Figure 2: Illustration of the space discretization.

Approximation of Anisotropic Terms

$\phi_{\sigma,v} : \sigma \longrightarrow [0, 1]$: Affine coordinates on triangle σ w.r.t. $v \in V(\sigma)$.

For each edge $[v, w]$ of some $\sigma \in \Sigma_m$:

$$\Sigma_{m,v,w} := \sigma \in \Sigma_m : \{v, w\} \subseteq V(\sigma) .$$

Letting

$$\Sigma_{\gamma_{m,v,w}} := \sigma \in \Sigma_{m,v,w} : \lambda_1(H_{v,w,\sigma} \cap \gamma_{m,v,w}) \neq 0 ,$$

decompose $\gamma_{m,v,w}$:

$$\gamma_{m,v,w} = \bigcup_{\sigma \in \Sigma_{\gamma_{m,v,w}}} \bar{\sigma} \cap \gamma_{m,v,w} .$$

Approximation:

$$(A_m \nabla \theta)|_{\sigma} \bullet \mathbf{n}_{\omega_v}|_{\gamma_{m,v,w}} \approx \sum_{\tilde{v} \in V(\sigma)} \theta(\tilde{v}) (A_m \nabla \phi_{\sigma,\tilde{v}}) \bullet \frac{w - v}{\|w - v\|_2} .$$

Finite Volume Scheme

Find $(\theta_v)_{v \in V}$ satisfying:

$$\theta_v = \theta_{\text{Dir}}(v) \quad \text{for each } v \in V_{\text{Dir}},$$

$$\begin{aligned} 0 = & \sum_{m \in M} \xi_m \theta_v - \theta_{\text{ext}}(v) \lambda_1(\partial\omega_{m,v} \cap \Gamma_{\text{Rob}}) \\ & - \sum_{m \in M} \sum_{\sigma \in \Sigma_{\gamma_{m,v,w}}} \frac{1}{2} \kappa_{\text{iso}}^m(\theta_v) + \kappa_{\text{iso}}^m(\theta_w) \\ & \sum_{\tilde{v} \in V(\sigma)} \theta_{\tilde{v}} (A_m \nabla \phi_{\sigma,\tilde{v}}) \bullet \frac{w - v}{\|w - v\|_2} \lambda_1(H_{v,w,\sigma} \cap \gamma_{m,v,w}) \\ & - \sum_{m \in M} f_{m,v} \lambda_2(\omega_{m,v}) \quad \text{for each } v \in V_{\neg\text{Dir}} = V \setminus V_{\text{Dir}} . \end{aligned}$$

Comparison with Closed-Form Solution

Axisymmetric Single-Material Domain

$$\Omega = \{(r, z) : 0 < r < 0.2, -0.2 < z < 0.2\}:$$

$$-\frac{1}{r} \frac{\partial}{\partial r} r \alpha_r \frac{\partial \theta}{\partial r} - \frac{\partial}{\partial z} \alpha_z \frac{\partial \theta}{\partial z} = 0 \quad \text{in } \Omega, \quad (1a)$$

$$\theta_{\text{Dir}}(r, z) := \frac{1}{2} \frac{1}{\alpha_r} r^2 - \frac{1}{\alpha_z} z^2 \quad \text{on } \partial\Omega. \quad (1b)$$

$$\text{Solution: } \theta(r, z) = \frac{1}{2} \frac{1}{\alpha_r} r^2 - \frac{1}{\alpha_z} z^2 \quad \text{on } \bar{\Omega}.$$

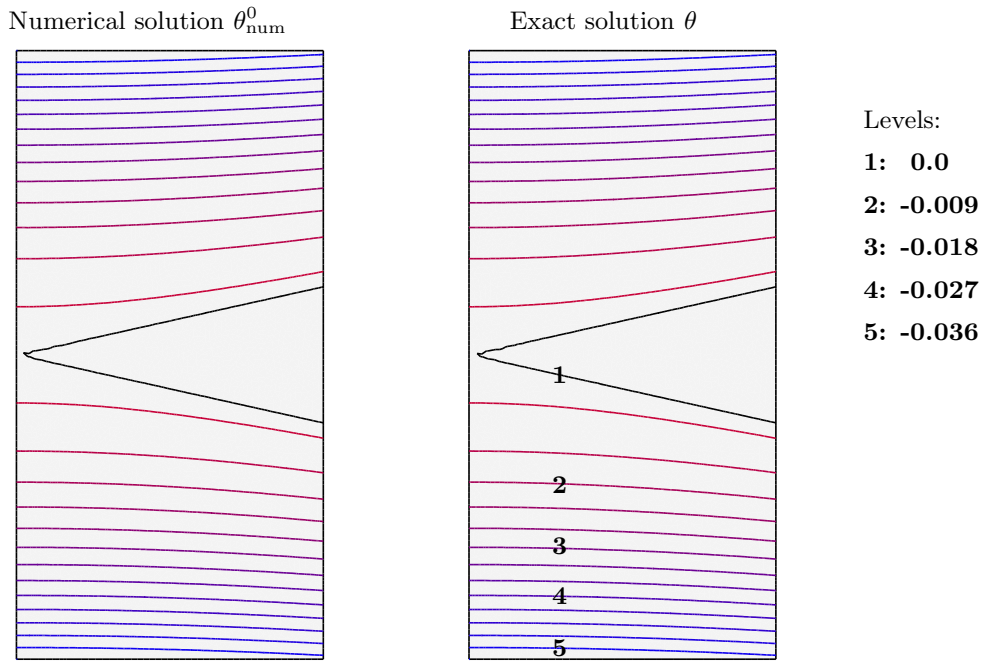


Figure 3: Solution of (1): Numerical θ_{num}^0 , 3117 triangles (left); exact θ (right). Isolevel difference: 0.003.

$$\text{Discrete } L_1\text{-error: } \epsilon_{L_1}^l := \sum_{v \in V^l} \text{vol}(\omega_v) |\theta_{\text{num}}^l(v) - \theta(v)|,$$

$v \in V^l$: vertices, $\text{vol}(\omega_v)$: r -weighted area of Voronoi cell.

Numerical convergence rate:

$$\rho_{L_1}^l := (\ln(\epsilon_{L_1}^l) - \ln(\epsilon_{L_1}^{l-1})) / (\ln(h^l) - \ln(h^{l-1})),$$

h^l : upper bound for triangle area of level l .

Axisymmetric Multi-Material Domain

$$\Omega_1 = \{(r, z) : 0 < r < r_0, 0 < z < z_0\},$$

$$\Omega_2 = \{(r, z) : r_0 < r < r_{\max}, 0 < z < z_0\},$$

$$\Omega_3 = \{(r, z) : 0 < r < r_0, z_0 < z < z_{\max}\},$$

$$\Omega_4 = \{(r, z) : r_0 < r < r_{\max}, z_0 < z < z_{\max}\},$$

$$-\frac{1}{r} \frac{\partial}{\partial r} \left(r \alpha_{m,r} \frac{\partial \theta}{\partial r} \right) - \frac{\partial}{\partial z} \left(\alpha_{m,z} \frac{\partial \theta}{\partial z} \right) = f_m \quad \text{in } \Omega_m, \quad (2a)$$

$$\left(\begin{pmatrix} \alpha_{m,r} & 0 \\ 0 & \alpha_{m,z} \end{pmatrix} \nabla \theta|_{\overline{\Omega}_m} \right) \bullet \mathbf{n}_m \quad (2b)$$

$$= \left(\begin{pmatrix} \alpha_{\tilde{m},r} & 0 \\ 0 & \alpha_{\tilde{m},z} \end{pmatrix} \nabla \theta|_{\overline{\Omega}_{\tilde{m}}} \right) \bullet \mathbf{n}_m \quad \text{on } \partial\Omega_m \cap \partial\Omega_{\tilde{m}},$$

$$\theta_{\text{Dir},m}(r, z) := a_m r^2 + b_m z^2 + c_m \quad \text{on } \partial\Omega \cap \partial\Omega_m. \quad (2c)$$

Solution:

$$\theta(r, z) := a_m r^2 + b_m z^2 + c_m \quad \text{on } \overline{\Omega}_m,$$

$$\theta_{\text{Dir},m}(r, z) := a_m r^2 + b_m z^2 + c_m \quad \text{on } \partial\Omega \cap \partial\Omega_m,$$

where $r_0 = z_0 = 0.1$, $r_{\max} = z_{\max} = 0.2$,

$$\alpha_{1,r} = 2, \quad \alpha_{2,r} = 1, \quad \alpha_{3,r} = 4, \quad \alpha_{4,r} = 2,$$

$$\alpha_{1,z} = 1, \quad \alpha_{2,z} = 2, \quad \alpha_{3,z} = 3, \quad \alpha_{4,z} = 6,$$

$$a_1 = 1, \quad a_2 = 2, \quad a_3 = 1, \quad a_4 = 2,$$

$$b_1 = 1, \quad b_2 = 1, \quad b_3 = 1/3, \quad b_4 = 1/3,$$

$$c_1 = 0, \quad c_2 = -1/100, \quad c_3 = 2/300, \quad c_4 = 1/300,$$

$$f_1 = -10, \quad f_2 = -12.0, \quad f_3 = -18.0, \quad f_4 = -20.0$$

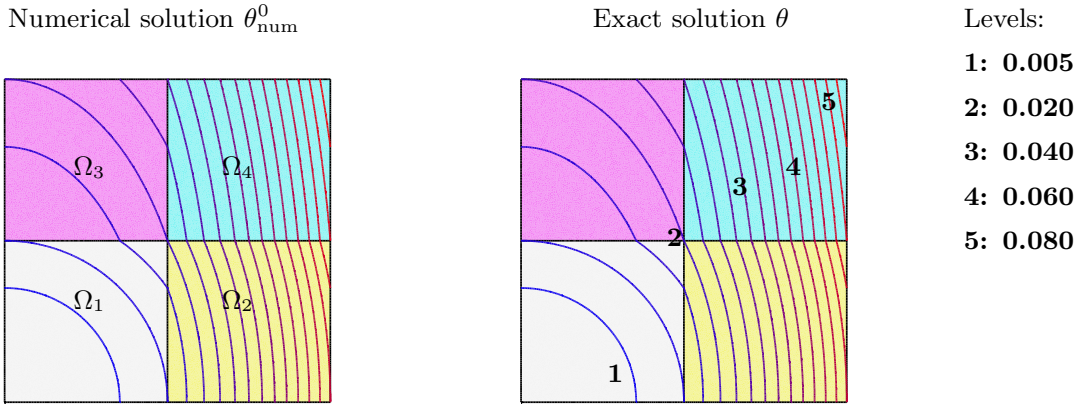


Figure 4: Solution of (2): Numerical θ_{num}^0 , 3117 triangles (left); exact θ (right). Isolevel difference: 0.005.

level l	number of triangles	max area h^l	L_1 -error $\epsilon_{L_1}^l$	numerical convergence rate $\rho_{L_1}^l$
0	3117	$4.0 \cdot 10^{-5}$	$5.013 \cdot 10^{-8}$	
1	12446	$1.0 \cdot 10^{-5}$	$1.260 \cdot 10^{-8}$	0.996125
2	49669	$2.5 \cdot 10^{-6}$	$3.244 \cdot 10^{-9}$	0.978789
3	198212	$6.25 \cdot 10^{-7}$	$8.2815 \cdot 10^{-10}$	0.984905
4	795195	$1.5625 \cdot 10^{-7}$	$2.0891 \cdot 10^{-10}$	0.993505

Table 1: L_1 -error and numerical convergence rate for the numerical solution of (1) with anisotropy $(\alpha_r, \alpha_z) = (10, 1)$.

level l	number of triangles	max area h^l	L_1 -error $\epsilon_{L_1}^l$	numerical convergence rate $\rho_{L_1}^l$
0	1557	$4.0 \cdot 10^{-5}$	$2.1325 \cdot 10^{-6}$	
1	6148	$1.0 \cdot 10^{-5}$	$1.0669 \cdot 10^{-6}$	0.49956
2	24813	$2.5 \cdot 10^{-6}$	$5.259 \cdot 10^{-7}$	0.510282
3	99428	$6.25 \cdot 10^{-7}$	$2.638 \cdot 10^{-7}$	0.497672
4	398130	$1.5625 \cdot 10^{-7}$	$1.3362 \cdot 10^{-7}$	0.490654

Table 2: L_1 -error and numerical convergence rate for the numerical solution of (2).

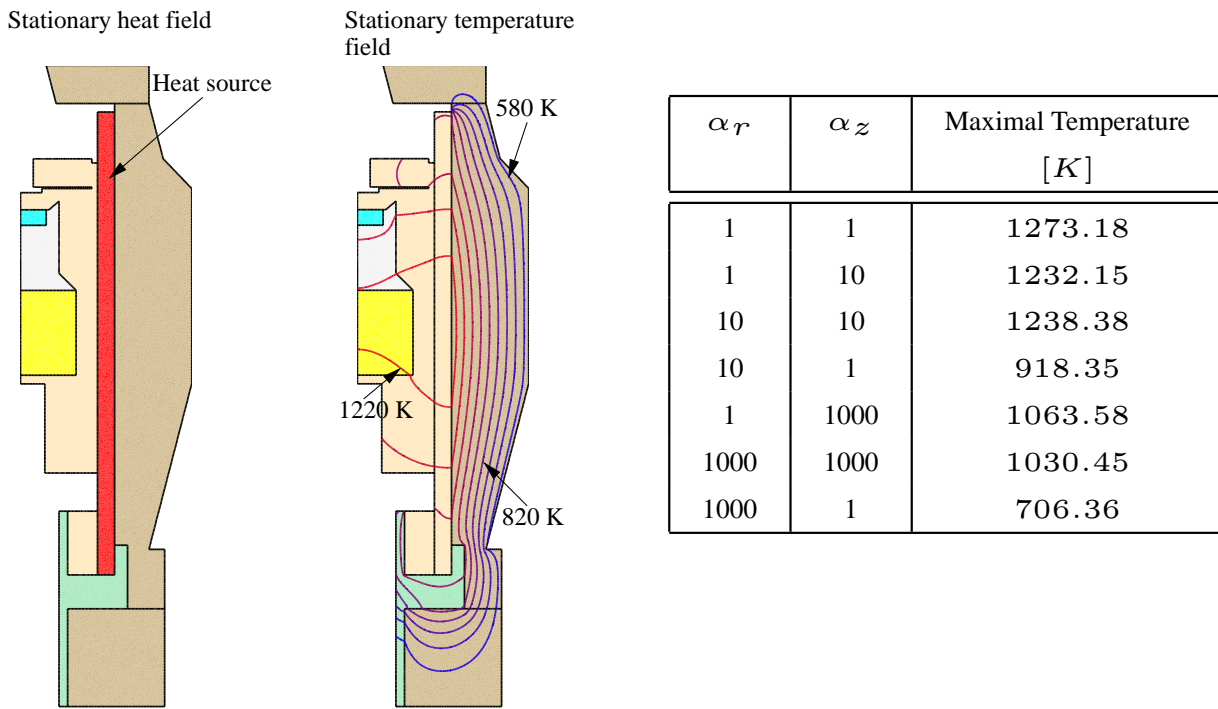


Figure 5: Left: Domain of heat sources highlighted. Right: T -field for isotropic insulation, i.e. $\alpha_r = \alpha_z = 1$.

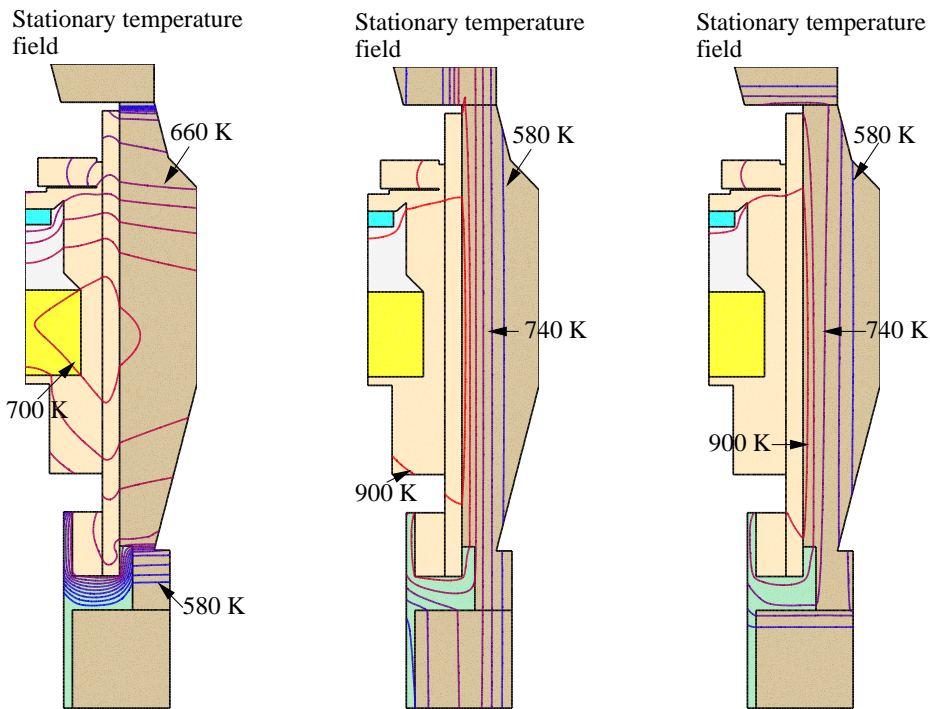


Figure 6: T -field for anisotropic insulation with $\alpha_r = 1000$ (left), with $\alpha_r = 1000$ for sides, $\alpha_z = 1000$ for top and bottom (middle), $\alpha_z = 1000$ (right).

Publications

- P. PHILIP: *Transient Numerical Simulation of Sublimation Growth of SiC Bulk Single Crystals. Modeling, Finite Volume Method, Results*, Thesis, Department of Mathematics, Humboldt University of Berlin, Germany, 2003 Report No. 22, Weierstrass Institute for Applied Analysis and Stochastics, Berlin.
- J. GEISER, O. KLEIN, P. PHILIP: *Numerical simulation of heat transfer in materials with anisotropic thermal conductivity: A finite volume scheme to handle complex geometries*. In preparation.
- J. GEISER, O. KLEIN, P. PHILIP: *Influence of anisotropic thermal conductivity in the apparatus insulation for sublimation growth of SiC: Numerical investigation of heat transfer*. In preparation.

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